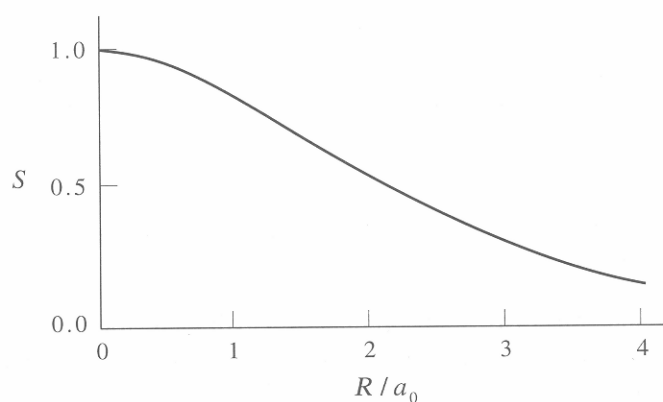


**FIGURE 9.4**  
The overlap of the 1s orbitals centered on hydrogen nuclei located at A and B, a distance  $R$  apart.

For two 1s orbitals, it can be shown that their overlap integral is

$$S(R) = e^{-R} \left( 1 + R + \frac{R^2}{3} \right) \quad (9.11)$$



**FIGURE 9.5**  
The overlap integral  $S(R)$ , Equation 9.11, for two hydrogen atom 1s orbitals plotted versus the internuclear separation in atomic units.

Evaluate numerator:

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Chap. 11 (p. 1)

$$\begin{aligned}\int \Psi_+^* \hat{H} \Psi_+ d\tau &= \int \Psi_+^* \left[ -\frac{\hbar^2}{2m} \nabla^2 + \frac{e^2}{4\pi\epsilon_0} \left( -\frac{1}{r_A} - \frac{1}{r_B} + \frac{1}{R} \right) \right] \Psi_+ d\tau \\ &= \int \Psi_+^* \left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{r_A} + \frac{1}{r_B} \right) \right] \Psi_+ d\tau \\ &\quad + \int \Psi_+^* \left[ \frac{e^2}{4\pi\epsilon_0 R} \right] \Psi_+ d\tau \\ &\equiv \textcircled{\text{I}} + \textcircled{\text{II}}\end{aligned}$$

$$\text{where } \textcircled{\text{II}} = \frac{e^2}{4\pi\epsilon_0 R} \int \Psi_+^* \Psi_+ d\tau = \frac{e^2}{4\pi\epsilon_0 R} 2N^2 [1 + S(R)]$$

↑ not integrating over R

Evaluate  $\textcircled{\text{I}}$  by substituting in LCAO expression:

$$\textcircled{\text{I}} = \int N(A+B) \left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{r_A} + \frac{1}{r_B} \right) \right] N(A+B) d\tau$$

Don't multiply out to get 12 terms!

$$\begin{aligned}\textcircled{\text{I}} &= N^2 \int (A+B) \left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{r_A} + \frac{1}{r_B} \right) \right] A d\tau \\ &\quad + N^2 \int (A+B) \left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{r_A} + \frac{1}{r_B} \right) \right] B d\tau\end{aligned}$$

(p.2)

and  $\left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r_A} \right] A = E(1s) A$  (eigenfunction!)

$\uparrow$   $e^-$  motion  
 (no nuclear motion)

$\uparrow$  distance of  $e^-$  in  $\Psi_{1s_A}$  from nucleus A

and  $\left[ -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r_B} \right] B = E(1s) B$  (eigenfunction!)

so  $\textcircled{\text{I}} = N^2 \int (A+B) \left[ E(1s) - \frac{e^2}{4\pi\epsilon_0 r_B} \right] A d\tau$   
 $+ N^2 \int (A+B) \left[ E(1s) - \frac{e^2}{4\pi\epsilon_0 r_A} \right] B d\tau$

and  $\int (A+B) E(1s) A d\tau = E(1s) \left[ \int A^2 d\tau + \int B A d\tau \right]$   
 $= E(1s) [1 + S(R)]$   
 $= \int (A+B) E(1s) B d\tau$

so  $\textcircled{\text{I}} = N^2 E(1s) [1 + S(R)] - N^2 \int (A+B) \left[ \frac{e^2}{4\pi\epsilon_0 r_B} \right] A d\tau$   
 $+ N^2 E(1s) [1 + S(R)] - N^2 \int (A+B) \left[ \frac{e^2}{4\pi\epsilon_0 r_A} \right] B d\tau$

We will not calculate, but we will label and interpret, the remaining integrals:

(I'll follow Atkins' notation on p.371)

(p. 3)

$$\text{Let } j \equiv \int A \left[ \frac{e^2}{4\pi\epsilon_0 r_B} \right] A d\tau = \frac{e^2}{4\pi\epsilon_0} \int \frac{A^2}{r_B} d\tau$$

This measures the interaction of the probability density of the left orbital ( $\psi_{1s_A}$ ) with the right nucleus ( $N_B$ )

By symmetry, must also be true that

$$j = \int B \left[ \frac{e^2}{4\pi\epsilon_0 r_A} \right] B d\tau = \frac{e^2}{4\pi\epsilon_0} \int \frac{B^2}{r_A} d\tau$$


---

$$\text{Let } k \equiv \int B \left[ \frac{e^2}{4\pi\epsilon_0 r_B} \right] A d\tau = \frac{e^2}{4\pi\epsilon_0} \int \frac{AB}{r_B} d\tau$$

This measures the interaction of the overlap density with the right nucleus.

Again, by symmetry, must also be true that

$$k = \int A \left[ \frac{e^2}{4\pi\epsilon_0 r_A} \right] B d\tau = \frac{e^2}{4\pi\epsilon_0} \int \frac{AB}{r_A} d\tau$$


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$$\text{so } \textcircled{I} = 2N^2 E(1s) [1 + S(R)] - N^2(j+k) - N^2(k+j)$$

$$\textcircled{I} = 2N^2 E(1s) [1 + S(R)] - 2N^2(j+k)$$